

# Critical behavior of the chain-generating function of self-avoiding walks on the Sierpinski gasket family: The Euclidean limit

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We study self-avoiding walks (SAW's) on the generalized Sierpinski gasket family of fractals. Each fractal can be labeled by an integer  $b$  ( $2 \leq b \leq \infty$ ), so that the fractal and spectral dimensions tend to the Euclidean value 2 when  $b \rightarrow \infty$ . By using an exact enumeration technique to obtain the series expansion for the chain-generating function of SAW's on these lattices, we calculate the associated critical exponent  $\gamma_b$  for  $2 \leq b \leq 100$ . The large- $b$  behavior of  $\gamma_b$  is the first numerical result consistent with the asymptotic convergence toward the Euclidean value  $\gamma_E$ . We also give an analytic argument supporting the assumption that  $\lim_{b \rightarrow \infty} \gamma_b \rightarrow \gamma_E$ . [S1063-651X(98)14008-4]

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## I. INTRODUCTION

Recently, increased attention has been focused on the limiting behavior of the critical properties of statistical systems on fractals when underlying fractal geometrical parameters such as the fractal ( $D_F$ ) or spectral ( $D_S$ ) dimensions approach the Euclidean values. This convergence is not ensured, because in order to obtain the critical properties of any statistical system on a fractal family (labeled by  $b$ ), one should analyze the results in the thermodynamic limit ( $N \rightarrow \infty$ ) for each  $b$ , while the critical properties on the limiting Euclidean lattice is obtained when the geometrical limit ( $b \rightarrow \infty$ ) is taken *before* the thermodynamic limit.

In this work we study the critical behavior of self-avoiding walks (SAW's) on a family of regular fractals embedded in the two-dimensional Euclidean space, the generalized Sierpinski gaskets (GSG's). Each member of the fractal family is labeled by an integer  $b \geq 2$ , and can be obtained as the result of an infinite iterative process in which a triangular structure is enlarged  $b$  times and a generator is reproduced in  $b(b+1)/2$  smallest triangles of the enlarged structure. The generator is the initial structure (see, for instance, Fig. 1 of Ref. [1]). For this fractal family, both  $D_F$  and  $D_S$  tends to the Euclidean value 2 when  $b \rightarrow \infty$ .

We present results for  $b \leq 100$  based on the series expansion method. The series expansion technique gives the most reliable results for the Euclidean lattices. This suggests the importance of extending it for fractal lattices.

We consider the chain-generating function for SAW's on a particular fractal,

$$C_b(x) = \sum_{n=1}^{\infty} c_n(b) x^n, \quad (1)$$

where  $c_n(b)$  is the number of distinct  $n$ -step SAW's per number of sites of the lattice and  $x$ , the fugacity, is the weight factor for each step. Near a critical fugacity  $x_c$ ,

$$C_b(x) \sim (x - x_c)^{-\gamma_b}, \quad (2)$$

where  $\gamma_b$  is the critical exponent and  $\mu_b \equiv (x_c)^{-1}$  is the connective constant.

The critical behavior of SAW's on the GSG family has recently been studied by series expansion method [1]. The critical fugacity  $\mu_b$  was numerically estimated, and it was found that  $\mu_b$  tends to the triangular value  $\mu_T$  when  $b \rightarrow \infty$ .

However, the asymptotic behavior of critical exponents is still a controversial issue. SAW's were studied on other fractal families which are considered to belong to the same universality class as the generalized Sierpinski gasket family studied here. Numerical results for the critical exponents  $\gamma_b$  and  $\nu_b$  were obtained via the Monte Carlo renormalization group (MCRG) for  $b \leq 80$  [2]. Although the range of  $b$  was not sufficiently large to allow a numerical estimate of the asymptotic behavior, the results of  $\gamma_b$  and  $\nu_b$  exhibit a monotonic behavior with  $b$  that depart from the respective Euclidean values ( $\gamma_E = 1.34$  and  $\nu_E = 0.75$ ) as  $b$  increases. On the other hand, finite-size scaling (FSS) arguments [3]

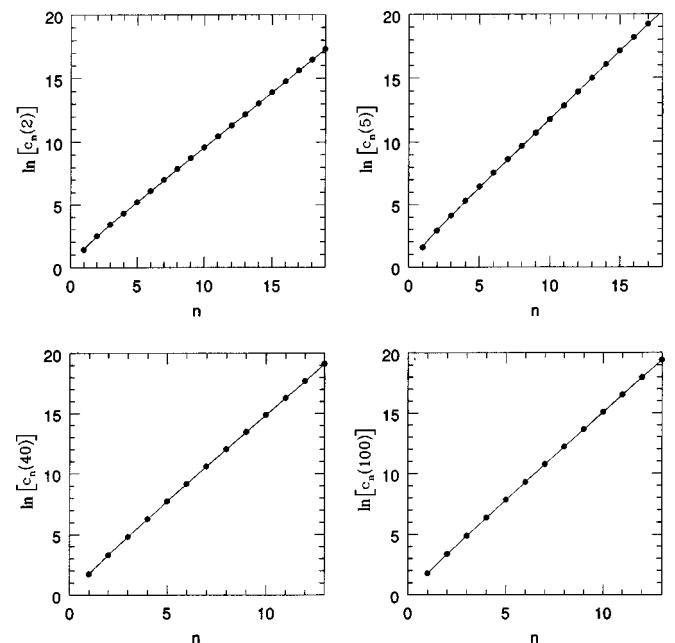


FIG. 1. Plot of  $\ln c_n(b)$  vs  $n$  for  $b=2, 5, 40$ , and  $100$ , respectively. For comparison, we plot the adjusted curve  $[\ln A_b + n \ln \mu_b + (\gamma_b - 1) \ln n]$  vs  $n$ .

$$\lim_{b \rightarrow \infty} X_b = \lim_{b \rightarrow \infty} U^{-1} Y_b = U^{-1} \lim_{b \rightarrow \infty} Y_b, \quad (13)$$

where  $U$  is  $b$  independent.

In Ref. [1], it was obtained that  $\lim_{b \rightarrow \infty} c_n(b) = c_n(T)$ , the density of SAW's of the underlying triangular lattice. Using Eqs. (7) and (8), then  $\lim_{b \rightarrow \infty} Y_b = Y_T$ , the corresponding value of the triangular lattice. Finally, from Eq. (13),

$$\lim_{b \rightarrow \infty} X_b = U^{-1} Y_T = X_T, \quad (14)$$

with  $X_T$  given by Eq. (8) with triangular parameters  $R_T$ ,  $S_T$ , and  $T_T$ .

Consequently, from Eqs. (5) and (14),

$$\lim_{b \rightarrow \infty} \mu_b = \mu_T \quad \text{and} \quad \lim_{b \rightarrow \infty} \gamma_b = \gamma_E.$$

#### IV. DISCUSSION

The convergence of critical exponents on fractals to those on uniform integer dimensional lattices is quite subtle, as explained in the text. In this work, we present results for the connective constant  $\mu_b$  and for the critical exponent  $\gamma_b$  of SAW's on a family of fractals that approaches the triangular lattice asymptotically as  $b \rightarrow \infty$ .

The SAW statistic is evaluated directly on the GSG family. Previous results in the literature regarding critical exponents for these lattices were obtained from other fractal families which are supposed to belong to the same universality class. For this reason they were not able to provide estimates for  $\mu_b$  which is a nonuniversal parameter.

We present analytic arguments supporting the convergence  $\lim_{b \rightarrow \infty} \mu_b = \mu_T$  and  $\lim_{b \rightarrow \infty} \gamma_b = \gamma_E$ . Our numerical estimates for  $\mu_b$  clearly shows that  $\lim_{b \rightarrow \infty} \mu_b = \mu_T$ , the connective constant of the underlying triangular lattice. From this, one can expect the analogous numerical convergence of the critical exponent.

Although the numerical estimates of  $\gamma_b$  were obtained for a large range of  $b$  ( $b \leq 100$ ), it was not sufficient to establish the numerical convergence of  $\gamma_b$  toward  $\gamma_E$ . That means that we have not reached the asymptotic regime, which would occur for larger  $b$ . Nevertheless, our values deviate by

at most 10% from  $\gamma_E$ , displaying a behavior in accordance with our analytic prediction.

In fact, to our knowledge this is the first numerical result showing this trend. Previous numerical findings[2], based on MCRG simulations of  $5 \times 10^6$  walks for  $b \leq 80$ , provide  $\gamma_b$  estimates that departs from  $\gamma_E$  as  $b$  increases. However, multiplying the number of sites of the fractal generator [ $N_S = (b+1)(b+2)/2$ ]—where the simulations were performed by our exact results for the density of  $n$ -step SAW's  $c_n(b)$ , shown in Table I, one finds that the number of  $n$ -step SAW's grows as  $b^2 c_n(b)$  which is very large compared with the number of Monte Carlo realizations, especially for large  $b$ . This could explain the disagreement between the MCGR results and the present work for large  $b$ .

From the MCGR data is not possible to obtain any limiting value for  $\gamma_b$  when  $b \rightarrow \infty$ . Although the authors argue that their results would be consistent with the limiting value  $\gamma_{FS} \approx 4.15$  provided by a FSS hypothesis [3], the largest value found in the simulations was  $\gamma \approx 2.2$ , which means a 50% deviate from the conjectured asymptotic value. In addition, the MCRG results for  $\nu_b$  seems inconsistent with the limiting value  $\nu_{FS} = \nu_E$  provided by the same FSS hypothesis. These controversial results call for additional studies.

Our numerical estimates of  $\mu_b$  and  $\gamma_b$  rely upon series expansions that are exact order by order. Each term of order  $n$  is obtained from an exact counting of  $c_n(b)$  for each infinite fractal lattice labeled by  $b$ . This means that the SAW statistics is calculated taking into account the existence of lacunas of all length scales, capturing the full geometry, in contrast with MC results that suffer from finite-size effects.

Finally, the series expansion method gives the most reliable results for Euclidean lattices. This gives confidence in the results based on the extension of this method for fractals.

The finite-size scaling predictions are based on a hypothesis regarding the dependence of critical quantities on  $b$  that has not been proved so far for the fractals families studied here. On the other hand, our analytic arguments are performed on a firm mathematical basis.

The results presented here leads to the conclusion that the critical behavior of SAW's on the SGS family of fractals shows a uniform convergence to the Euclidean behavior as  $b \rightarrow \infty$ .

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